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The Laplace Decomposition Method for Solving Nonlinear Conformable Fractional Evolution Equations

Ahmed ANBER and Zoubir DAHMANI

Department of Mathematics University of USTO, Oran, 31000, Algeria

Department of Mathematics, University of Blida 1, Blida 09000, Algeria
e-mail: ah.anber@gmail.com, zzdahmani@yahoo.fr

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Abstract

In the present paper, we discuss the application of the Laplace decomposition method (LDM) to find approximate solutions for some classes of time and space-conformable fractional evolution equations. Some illustrative examples are presented showing that the LDM is an efficient method for finding solutions of nonlinear conformable fractional problems.

Keywords: *Conformable derivative, Laplace Decomposition Method, Fractional differential equation.*

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1 Introduction

Some applications in sciences are modeled by nonlinear fractional differential equations and it is very difficult to find solutions to these equations. For this reason, there are several methods that can be applied to find approximate solutions [5, 6, 10, 11, 28]. In particular, one of these methods, is the Laplace decomposition method (LDM), for more information and also for some applications, we cite the research papers [7, 8, 13, 14, 29, 31, 34]. The principle of the LDM method is in combining the Laplace transform method [12, 21, 23] and the Adomian decomposition method [2, 4, 22, 26].

In this paper, we apply the LDM method to solve some nonlinear conformable

fractional differential equations of the form:

$$T_t^\beta u(x, t) + aT_x^\alpha u(x, t) + bu(x, t) + g(u(x, t)) = h(x, t), \quad (1.1)$$

with the initial conditions

$$u(x, 0) = f_1(x), \quad u_t(x, 0) = f_2(x), \quad (1.2)$$

where T_x^β, T_t^α are the conformable fractional derivatives in the sense of Khalil [18], with $1 < \alpha, \beta \leq 2$, a, b are constants, $h(x, t)$ is the inhomogeneous part, $g(u(x, t))$ is a nonlinear function of $u(x, t)$ and $f_1(x), f_2(x)$ are a given functions.

The paper is organized as follows: In the second section, some preliminaries related to conformable fractional derivative approach are recalled. In the third section, the Laplace decomposition method is discussed for solving some important conformable fractional evolution equations. In the fourth section, numerical examples are presented. Finally, some graphs of the obtained solutions are plotted, and a conclusion follows.

2 Conformable Fractional Concepts

In this section, we introduce some definitions and properties, see [1, 3, 9, 18, 20, 25, 32].

Definition 2.1 Let $f : (0, \infty) \rightarrow \mathbb{R}$. The conformable fractional derivative of order α is defined by

$$(T^\alpha f)(t) = \frac{\partial^\alpha f(t, x)}{\partial t^\alpha} = \lim_{\varepsilon \rightarrow 0} \left(\frac{f(t + \varepsilon t^{1-\alpha}) - f(t)}{\varepsilon} \right), \quad t > 0, \quad 0 < \alpha \leq 1. \quad (2.1)$$

Definition 2.2 The conformable fractional integral of a function $f : (0, \infty) \rightarrow \mathbb{R}$ of order α is defined as

$$(I^\alpha f)(t) = \int_0^t \tau^{\alpha-1} f(\tau) d\tau, \quad 0 < \alpha \leq 1. \quad (2.2)$$

In this paper, we need to recall the properties:

$$I^\alpha T^\alpha f(t) = f(t) - f(0) \quad (2.3)$$

and

$$(T^\alpha f)(t) = t^{1-\alpha} \frac{df(t)}{dt}. \quad (2.4)$$

Definition 2.3 Let $f : (0, \infty) \rightarrow \mathbb{R}$ be a real valued function and $0 < \alpha \leq 1$. Then the conformable fractional Laplace transform of f is defined as:

$$L_\alpha [f(x)](s) = {}_\alpha (s) = \int_0^\infty e^{-s(\frac{x^\alpha}{\alpha})} f(x) d_\alpha x = \int_0^\infty e^{-s(\frac{x^\alpha}{\alpha})} f(x) x^{\alpha-1} dx \quad (2.5)$$

The Laplace transform for the conformable fractional-order derivative is described as follows:

$$L_\alpha (T^\alpha f(x)) = sL_\alpha [f(x)] - f(0) \quad (2.6)$$

The relation between the usual and the fractional Laplace transforms is given below.

Theorem 2.4 [1] Let $f : (0, \infty) \rightarrow \mathbb{R}$ be a real valued function such that $L_\alpha [f(x)](s) = {}_\alpha (s)$ exists. Then:

$${}_\alpha (s) = L \left[f \left((\alpha x)^{\frac{1}{\alpha}} \right) \right] (s),$$

where $L[g(x)](s) = \int_0^\infty e^{-st} f(x) dx$

It is easy to show that:

Theorem 2.5 If $L_\alpha [f(x)](s) = {}_\alpha (s)$ exists. Then:

$$L_\alpha [c](s) = \frac{c}{s}$$

$$L_\alpha [x^p](s) = \alpha^{\frac{p}{\alpha}} \frac{\Gamma(1+\frac{p}{\alpha})}{s^{1+\frac{p}{\alpha}}}, \quad c \text{ and } p \text{ are arbitrary constants.}$$

3 Main Steps of LDM Method

In this section, we present the Laplace Decomposition methods in the case of Khalil fractional theory. For more details on the two methods, one can consult [15, 16, 19, 27, 30].

Firstly, Laplace transform is applied on both sides of the equation (1.1). Then, we get

$$L_\beta \left[T_t^\beta u(x, t) \right] + aL_\beta [T_x^\alpha u(x, t)] + bL_\beta [u(x, t)] + L_\beta [g(u(x, t))] = L_\beta [h(x, t)], \quad (3.1)$$

hence,

$$\begin{aligned} s^\beta u(x, s) - s^{\beta-1} u(x, 0) - s^{\beta-2} u_t(x, 0) &= L_\beta [h(x, t)] - aL_\beta [T_x^\alpha u(x, t)] \\ &\quad - bL_\beta [u(x, t)] - L_\beta [g(u(x, t))], \end{aligned} \quad (3.2)$$

equ. (3.2) can be written as

$$\begin{aligned} u(x, s) = & \frac{1}{s}u(x, 0) + \frac{1}{s^2}u_t(x, 0) + \frac{1}{s^\beta}L_\beta[h(x, t)] - \frac{a}{s^\beta}L_\beta[T_x^\alpha u(x, t)] \\ & - \frac{b}{s^\beta}L_\beta[u(x, t)] - \frac{1}{s^\beta}L_\beta[g(u(x, t))], \end{aligned} \quad (3.3)$$

Now, we use the fact that the Adomian decomposition method assumes that the function $u(x, t)$ can be decomposed into an infinite series

$$u(x, t) = \sum_{n=0}^{\infty} u_n(x, t) \quad (3.4)$$

It also assumes that

$$g(u(x, t)) = \sum_{n=0}^{\infty} A_n, \quad (3.5)$$

where A_n are Adomian polynomials given by

$$A_n = \frac{1}{n!} \frac{d^n}{d\lambda^n} \left[g \left(\sum_{i=0}^{\infty} \lambda^i u_i \right) \right]_{\lambda=0}, \quad n = 0, 1, 2, \dots \quad (3.6)$$

where λ is a parameter.

So, we have

$$\begin{aligned} A_0 &= \frac{1}{0!} \frac{d^0}{d\lambda^0} \left[g \left(\sum_{i=0}^{\infty} \lambda^i u_i \right) \right]_{\lambda=0} = g(u_0) \\ A_1 &= \frac{1}{1!} \frac{d^1}{d\lambda^1} \left[g \left(\sum_{i=0}^{\infty} \lambda^i u_i \right) \right]_{\lambda=0} = u_1 g'(u_0) \\ A_2 &= \frac{1}{2!} \frac{d^2}{d\lambda^2} \left[g \left(\sum_{i=0}^{\infty} \lambda^i u_i \right) \right]_{\lambda=0} = \frac{1}{2!} (u_1)^2 g''(u_0) + u_2 g'(u_0) \\ &\vdots \\ &\vdots \\ &\vdots \end{aligned}$$

Substitution Equation (3.4), Equation (3.5) and Equation (3.6) in Equation (3.3), we have

$$\begin{aligned} \sum_{n=0}^{\infty} u_n(x, s) = & \frac{1}{s}f_1(x) + \frac{1}{s^2}f_2(x) + \frac{1}{s^\beta}L_\beta[h(x, t)] - \frac{a}{s^\beta}L_\beta \left[T_x^\alpha \left(\sum_{n=0}^{\infty} u_n(x, t) \right) \right] \\ & - \frac{b}{s^\beta}L_\beta \left[\sum_{n=0}^{\infty} u_n(x, t) \right] - \frac{1}{s^\beta}L_\beta \left[\sum_{n=0}^{\infty} A_n \right], \end{aligned} \quad (3.7)$$

Hence,

$$\begin{aligned}
u_0(x, s) &= \frac{1}{s} f_1(x) + \frac{1}{s^2} f_2(x) + \frac{1}{s^\beta} L_\beta [h(x, t)] = f_3(x, s) \\
u_1(x, s) &= -\frac{a}{s^\beta} L_\beta [T_x^\alpha (u_0(x, t))] - \frac{b}{s^\beta} L_\beta [u_0(x, t)] - \frac{1}{s^\beta} L_\beta [A_0] \\
&\cdot \\
&\cdot \\
&\cdot \\
u_{n+1}(x, s) &= -\frac{a}{s^\beta} L_\beta [T_x^\alpha (u_n(x, t))] - \frac{b}{s^\beta} L_\beta [u_n(x, t)] - \frac{1}{s^\beta} L_\beta [A_n]
\end{aligned} \tag{3.8}$$

When the inverse Laplace transform Equation (3.8) is applied, we get

$$\begin{aligned}
u_0(x, t) &= f_3(x, t) \\
u_{n+1}(x, t) &= -L^{-1} \left[\frac{a}{s^\beta} L_\beta (a T_x^{2\alpha} (u_n(x, t)) + b u_n(x, t) + A_n) \right]
\end{aligned} \tag{3.9}$$

4 Nonlinear Applications

As application, we begin by the following nonlinear example:

4.1 Example 1

We consider the nonlinear fractional equation

$$T_t^\beta u(x, t) - T_x^\alpha u(x, t) = \lambda u^\gamma(x, t), \quad 1 < \alpha, \beta \leq 2. \tag{4.1}$$

with the initial conditions:

$$u(x, 0) = g_1(x), \quad u_t(x, 0) = g_2(x). \tag{4.2}$$

We have

$$u(x, s) = \frac{1}{s} u(x, 0) + \frac{1}{s^2} u_t(x, 0) + \frac{1}{s^\beta} L_\beta [T_x^\alpha u(x, t)] + \frac{\lambda}{s^\beta} L_\beta [u^\gamma(x, t)] \tag{4.3}$$

By using the initial condition, the following recurrence relational is obtained

$$u(x, s) = \frac{g_1(x)}{s} + \frac{g_2(x)}{s^2} + \frac{1}{s^\beta} L_\beta (T_x^\alpha u(x, t) + \lambda u^\gamma(x, t)) \tag{4.4}$$

Applying the inverse Laplace transform, to both sides of Equation (4.4) yields:

$$u(x, t) = g_1(x) + t g_2(x) + L_\beta^{-1} [T_x^\alpha u(x, t) + \lambda u^\gamma(x, t)] \tag{4.5}$$

Case 1: For $\gamma = 1$

We have

$$\sum_{n=0}^{\infty} u_n(x, t) = g_1(x) + t g_2(x) + L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta \left[\sum_{n=0}^{\infty} T_x^\alpha u_n(x, t) + \lambda \sum_{n=0}^{\infty} u_n(x, t) \right] \right]$$

Hence,

$$\begin{aligned} u_0(x, t) &= g_1(x) + tg_2(x) \\ u_{n+1}(x, t) &= L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta \left[\sum_{n=0}^{\infty} T_x^\alpha u_n(x, t) + \lambda \sum_{n=0}^{\infty} u_n(x, t) \right] \right] \end{aligned} \quad (4.6)$$

Consequently

$$u_1(x, t) = L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta (T_x^\alpha u_0(x, t) + \lambda u_0(x, t)) \right] = (\lambda g_1(x) + T^\alpha(g_1(x))) \frac{t^\beta}{\Gamma(\beta+1)} + (\lambda g_2(x) + T^\alpha(g_2(x))) \frac{t^{\beta+1}}{\Gamma(\beta+2)}$$

$$\begin{aligned} u_2(x, t) &= L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta (T_x^\alpha u_1(x, t) + \lambda u_1(x, t)) \right] \\ &= L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta \left(\begin{array}{l} T_x^\alpha \left(\begin{array}{l} ((\lambda g_1(x) + T_x^\alpha(g_1(x))) \frac{t^\beta}{\Gamma(\beta+1)}) \\ + (\lambda g_2(x) + T_x^\alpha(g_2(x))) \frac{t^{\beta+1}}{\Gamma(\beta+2)} \end{array} \right) \\ + \lambda \left(\begin{array}{l} (\lambda g_1(x) + T_x^\alpha(g_1(x))) \frac{t^\beta}{\Gamma(\beta+1)} \\ + (\lambda g_2(x) + T_x^\alpha(g_2(x))) \frac{t^{\beta+1}}{\Gamma(\beta+2)} \end{array} \right) \end{array} \right) \right] \end{aligned}$$

$$\begin{aligned} &= [2\lambda T_x^\alpha(g_1(x)) + T_x^{2\alpha}(g_1(x)) + \lambda^2 g_1(x)] \frac{t^{2\beta}}{\Gamma(2\beta+1)} \\ &+ [2\lambda T_x^\alpha(g_2(x)) + T_x^{2\alpha}(g_2(x)) + \lambda^2 g_2(x)] \frac{t^{2\beta+1}}{\Gamma(2\beta+2)} \end{aligned}$$

$$\begin{aligned} u_3(x, t) &= L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta \left(\frac{\partial^\alpha u_2(x, t)}{\partial x^\alpha} + \lambda u_2(x, t) \right) \right] \\ &= L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta \left(\begin{array}{l} T_x^\alpha \left(\begin{array}{l} [2\lambda T_x^\alpha(g_1(x)) + T_x^{2\alpha}(g_1(x)) + \lambda^2 g_1(x)] \frac{t^{2\beta}}{\Gamma(2\beta+1)} \\ [2\lambda T_x^\alpha(g_2(x)) + T_x^{2\alpha}(g_2(x)) + \lambda^2 g_2(x)] \frac{t^{2\beta+1}}{\Gamma(2\beta+2)} \end{array} \right) \\ + \lambda \left(\begin{array}{l} [2\lambda T_x^\alpha(g_1(x)) + T_x^{2\alpha}(g_1(x)) + \lambda^2 g_1(x)] \frac{t^{2\beta}}{\Gamma(2\beta+1)} \\ [2\lambda T_x^\alpha(g_2(x)) + T_x^{2\alpha}(g_2(x)) + \lambda^2 g_2(x)] \frac{t^{2\beta+1}}{\Gamma(2\beta+2)} \end{array} \right) \end{array} \right) \right] \\ &= [3\lambda T_x^{2\alpha}(g_1(x)) + 3\lambda^2 T_x^\alpha(g_1(x)) + T_x^{3\alpha}(g_1(x)) + \lambda^3 g_1(x)] \frac{t^{3\beta}}{\Gamma(3\beta+1)} \\ &+ [3\lambda T_x^{2\alpha}(g_2(x)) + 3\lambda^2 T_x^\alpha(g_2(x)) + T_x^{3\alpha}(g_2(x)) + \lambda^3 g_2(x)] \frac{t^{3\beta+1}}{\Gamma(3\beta+2)} \end{aligned}$$

⋮
⋮
⋮

For $\lambda = 1$, $g_1(x) = 1 + \sin x$, $g_2(x) = 0$ and $\alpha = 2$, we obtain

$$u_0(x, t) = 1 + \sin x$$

$$u_1(x, t) = \frac{t^\beta}{\Gamma(\beta+1)}$$

⋮

⋮

⋮

$$u_n(x, t) = \frac{t^{n\beta}}{\Gamma(n\beta+1)}$$

Thus,

$$u(x, t) = 1 + \sin x + \frac{t^\beta}{\Gamma(\beta+1)} + \frac{t^{2\beta}}{\Gamma(2\beta+1)} + \dots + \frac{t^{n\beta}}{\Gamma(n\beta+1)} + \dots \quad (4.7)$$

To find the validity of the approximated solution, when $\beta = 2$, the exact solution is:

$$u(x, t) = \sin x + 1 + \frac{t^2}{\Gamma(3)} + \frac{t^4}{\Gamma(5)} + \dots + \frac{t^{2n}}{\Gamma(2n+1)} + \dots = \sin x + \cosh t. \quad (4.8)$$

Figure 1.1 is a graph of the solution (4.7), respectively for β equal to 1.25, 1.5, 1.75, and 2.

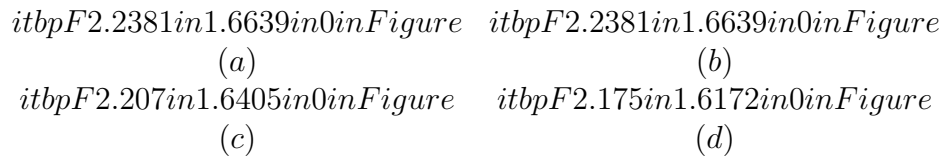


Fig1.1 Graph of Eq.(4.7) with different β values
 (a) $\beta = 1.25$, (b) $\beta = 1.5$, (c) $\beta = 1.75$, (d) $\beta = 2$,

Case 2: For $\gamma \neq 1$

We have

$$\sum_{n=0}^{\infty} u_n(x, t) = g_1(x) + tg_2(x) + L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta \left[\sum_{n=0}^{\infty} T_x^\alpha u_n(x, t) + \lambda \sum_{n=0}^{\infty} A_n \right] \right] \quad (4.9)$$

The few components of Adomian polynomials are given by

$$\begin{aligned} A_0 &= u_0^\gamma \\ A_1 &= \gamma u_1 u_0^{\gamma-1} \\ A_2 &= \gamma u_2 u_0^{\gamma-1} + \frac{\gamma(\gamma-1)}{2} u_1^2 u_0^{\gamma-2} \\ A_3 &= \gamma u_3 u_0^{\gamma-1} + \gamma(\gamma-1) u_1 u_2 u_0^{\gamma-2} + \frac{\gamma(\gamma-1)(\gamma-2)}{6} u_1^3 u_0^{\gamma-3} \end{aligned} \quad (4.10)$$

Therefore,

$$\begin{aligned} u_0(x, t) &= g_1(x) + tg_2(x) \\ u_{n+1}(x, t) &= L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta \left[\sum_{n=0}^{\infty} T_x^\alpha u_n(x, t) + \lambda \sum_{n=0}^{\infty} A_n \right] \right] \end{aligned} \quad (4.11)$$

Consequently

$$\begin{aligned}
 u_1(x, t) &= L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta (T_x^\alpha u_0(x, t) + \lambda A_0) \right] \\
 &= L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta (T_x^\alpha u_0 + \lambda u_0^\gamma) \right] \\
 u_2(x, t) &= L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta (T_x^\alpha u_1(x, t) + \lambda A_1) \right] \\
 &= L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta (T_x^\alpha u_1 + \lambda \gamma u_1 u_0^{\gamma-1}) \right] \\
 u_3(x, t) &= L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta (T_x^\alpha u_2(x, t) + \lambda A_2) \right] \\
 &= L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta \left(T_x^\alpha u_2 + \lambda \gamma u_2 u_0^{\gamma-1} + \frac{\lambda \gamma (\gamma-1)}{2} u_1^2 u_0^{\gamma-2} \right) \right] \\
 &\vdots \\
 &\vdots \\
 &\vdots
 \end{aligned}$$

For $\lambda = 1, \gamma = 2, g_1(x) = x^2, g_2(x) = 0$, we obtain

$$\begin{aligned}
 u_0(x, t) &= x^2 \\
 u_1(x, t) &= L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta (T_x^\alpha u_0(x, t) + A_0) \right] \\
 &= L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta (T_x^\alpha u_0(x, t) + u_0^2) \right] = L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta (2x^{2-\alpha} + x^4) \right] \\
 &= (2x^{2-\alpha} + x^4) \frac{t^\beta}{\Gamma(\beta+1)} \\
 u_2(x, t) &= L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta (T_x^\alpha u_1(x, t) + A_1) \right] \\
 &= L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta \left((2(2-\alpha)(1-\alpha)x^{2-2\alpha} + 16x^{4-\alpha} + 2x^6) \frac{t^\beta}{\Gamma(\beta+1)} \right) \right] \\
 &= (2(2-\alpha)(1-\alpha)x^{2-2\alpha} + 16x^{4-\alpha} + 2x^6) \frac{t^{2\beta}}{\Gamma(2\beta+1)} \\
 &\vdots \\
 &\vdots \\
 &\vdots
 \end{aligned} \tag{4.12}$$

Thus,

$$\begin{aligned}
 u(x, t) &= x^2 + (2x^{2-\alpha} + x^4) \frac{t^\beta}{\Gamma(\beta+1)} \\
 &+ (2(2-\alpha)(1-\alpha)x^{2-2\alpha} + 16x^{4-\alpha} + 2x^6) \frac{t^{2\beta}}{\Gamma(2\beta+1)} + \dots
 \end{aligned} \tag{4.13}$$

Figure 1.2 is a graph of the solution (4.13), for α, β equal to 1.25, 1.5, 1.75, and 2.

$$\begin{array}{cc}
 \textit{itbpF2.29in1.702in0inFigure} & \textit{itbpF2.3004in1.7097in0inFigure} \\
 (a) & (b) \\
 \textit{itbpF2.2174in1.6483in0inFigure} & \textit{itbpF2.2174in1.6483in0inFigure} \\
 (c) & (d)
 \end{array}$$

Fig1.2 Graph of Eq.(4.13) with different α, β values

(a) $\alpha = \beta = 1.25$, (b) $\alpha = \beta = 1.5$, (c) $\alpha = \beta = 1.75$, (d) $\alpha = \beta = 2$,

4.2 Example 2

We consider the nonlinear fractional equation

$$T_t^\beta u(x, t) - T_x^\alpha u(x, t) = u(x, t) + u^2(x, t), \quad 1 < \alpha, \beta \leq 2. \quad (4.14)$$

with the initial conditions:

$$u(x, 0) = 1 + x, \quad u_t(x, 0) = 0. \quad (4.15)$$

When the Laplace transform is applied to both sides of Equation (4.14), we can write

$$u(x, s) = \frac{1}{s}u(x, 0) + \frac{1}{s^2}u_t(x, 0) + \frac{1}{s^\beta}L_\beta [T_x^\alpha u(x, t)] + \frac{1}{s^\beta}L_\beta [u^2(x, t)] \frac{1}{s^\beta}L_\beta [u(x, t)] \quad (4.16)$$

Hence,

$$u(x, s) = \frac{x^2}{s} + \frac{1}{s^\beta}L_\beta (T_x^\alpha u(x, t) + u^2(x, s) + u(x, s)) \quad (4.17)$$

Applying the inverse Laplace transform, to both sides of Equation (4.17) yields:

$$u(x, t) = x^2 + L_\beta^{-1} \left[\frac{1}{s^\beta}L_\beta (T_x^\alpha u(x, t) + u^2(x, t) + u(x, t)) \right] \quad (4.18)$$

Hence,

$$\sum_{n=0}^{\infty} u_n(x, t) = x^2 + L_\beta^{-1} \left[\frac{1}{s^\beta}L_\beta \left[\sum_{n=0}^{\infty} T_x^\alpha u_n(x, t) + \sum_{n=0}^{\infty} B_n + u_n(x, t) \right] \right] \quad (4.19)$$

On the other hand, we have

$$\begin{aligned} B_0 &= u_0^2 \\ B_1 &= 2u_0u_1 \\ B_2 &= 2u_0u_2 + u_1^2 \\ B_3 &= 2u_0u_3 + 2u_1u_2 \end{aligned}$$

Then

$$\begin{aligned} u_0(x, t) &= 1 + x^2, \\ u_{n+1}(x, t) &= L_\beta^{-1} \left[\frac{1}{s^\beta}L_\beta \left[\sum_{n=0}^{\infty} T_x^\alpha u_n(x, t) + \sum_{n=0}^{\infty} B_n + u_n(x, t) \right] \right] \end{aligned} \quad (4.20)$$

Consequently

$$\begin{aligned} u_1(x, t) &= L_\beta^{-1} \left[\frac{1}{s^\beta}L_\beta (T_x^\alpha u_0(x, t) + B_0 + u_0(x, t)) \right] \\ &= L_\beta^{-1} \left[\frac{1}{s^\beta}L_\beta (T_x^\alpha u_0 + u_0^2 + u_0) \right] \\ &= L_\beta^{-1} \left[\frac{1}{s^\beta}L_\beta (2x^{2-\alpha} + x^4 + 2 + 3x^2) \right] \\ &= (2x^{2-\alpha} + x^4 + 3x^2 + 2) \frac{t^\beta}{\Gamma(\beta+1)} \end{aligned}$$

$$\begin{aligned}
u_2(x, t) &= L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta (T_x^\alpha u_1(x, t) + B_1 + u_1(x, t)) \right] \\
&= L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta (T_x^\alpha u_1 + 2u_0 u_1 + u_1) \right] \\
&= L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta \left(\begin{array}{c} T_x^\alpha \left((2x^{2-\alpha} + x^4 + 3x^2 + 2) \frac{t^\beta}{\Gamma(\beta+1)} \right) \\ + 2(1+x^2) \left((2x^{2-\alpha} + x^4 + 3x^2 + 2) \frac{t^\beta}{\Gamma(\beta+1)} \right) \\ + (2x^{2-\alpha} + x^4 + 3x^2 + 2) \frac{t^\beta}{\Gamma(\beta+1)} \end{array} \right) \right] \\
&= L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta \left(\begin{array}{c} T_x^\alpha \left((2x^{2-\alpha} + x^4 + 3x^2 + 2) \frac{t^\beta}{\Gamma(\beta+1)} \right) \\ + (6x^{2-\alpha} + 4x^{4-\alpha} + 2x^6 + 9x^4 + 13x^2 + 6) \frac{t^\beta}{\Gamma(\beta+1)} \end{array} \right) \right] \\
&= (2(2-\alpha)x^{2-2\alpha} + 8x^{4-\alpha} + 12x^{2-\alpha} + 2x^6 + 9x^4 + 13x^2 + 6) \frac{t^{2\beta}}{\Gamma(2\beta+1)} \\
&\vdots \\
&\vdots \\
&\vdots
\end{aligned}$$

Thus,

$$\begin{aligned}
u(x, t) &= 1 + x^2 + (2x^{2-\alpha} + x^4 + 3x^2 + 2) \frac{t^\beta}{\Gamma(\beta+1)} \\
&+ (2(2-\alpha)x^{2-2\alpha} + 8x^{4-\alpha} + 12x^{2-\alpha} + 2x^6 + 9x^4 + 13x^2 + 6) \frac{t^{2\beta}}{\Gamma(2\beta+1)} \dots
\end{aligned} \tag{4.21}$$

Figure 1.3 is a graph of the solution (4.21), for α, β equal to 1.25, 1.5, 1.75, and 2.

$$\begin{array}{cc}
\text{itbpF2.3627in1.7582in0inFigure} & \text{itbpF2.373in1.7651in0inFigure} \\
(a) & (b) \\
\text{itbpF2.4146in1.7953in0inFigure} & \text{itbpF2.4362in1.8118in0inFigure} \\
(c) & (d)
\end{array}$$

Fig1.3 Graph of Eq.(4.21) with different α, β values

(a) $\alpha = \beta = 1.25$, (b) $\alpha = \beta = 1.5$, (c) $\alpha = \beta = 1.75$, (d) $\alpha = \beta = 2$,

4.3 Example 3

Consider the nonlinear fractional wave equation [16, 17]

$$\begin{aligned}
T_t^\alpha u(x, t) + aT_x^\alpha u(x, t) + bu(t, x) + cu^3(t, x) &= 0, \\
0 < x < 1, t > 0, 1 < \alpha, \beta \leq 2.
\end{aligned} \tag{4.22}$$

with the initial conditions:

$$u(x, 0) = h_1(x), \quad u_t(x, 0) = h_2(x). \tag{4.23}$$

where a, b, c are constants.

We have

$$\begin{aligned} u_0(x, t) &= h_1(x) + t \times h_2(x) \\ u_{n+1}(x, t) &= -L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta \left[a \sum_{n=0}^{\infty} T_x^\alpha u_n(x, t) + b \sum_{n=0}^{\infty} u_n(x, t) + c \sum_{n=0}^{\infty} C_n \right] \right] \end{aligned} \quad (4.24)$$

where C_n are given by

$$C_n = \frac{1}{n!} \frac{d^n}{d\lambda^n} \left[\left(\sum_{i=0}^{\infty} \lambda^i u_i \right)^3 \right]_{\lambda=0}, \quad n = 0, 1, 2, \dots \quad (4.25)$$

Hence,

$$\begin{aligned} C_0 &= \frac{1}{0!} \frac{d^0}{d\lambda^0} \left[\left(\sum_{i=0}^{\infty} \lambda^i u_i \right)^3 \right]_{\lambda=0} = (u_0)^3 \\ C_1 &= \frac{1}{1!} \frac{d^1}{d\lambda^1} \left[\left(\sum_{i=0}^{\infty} \lambda^i u_i \right)^3 \right]_{\lambda=0} = 3u_1 (u_0)^2 \\ C_2 &= \frac{1}{2!} \frac{d^2}{d\lambda^2} \left[\left(\sum_{i=0}^{\infty} \lambda^i u_i \right)^3 \right]_{\lambda=0} = 3(u_1)^2 (u_0) + 3u_2 (u_0)^2 \\ C_3 &= \frac{1}{3!} \frac{d^3}{d\lambda^3} \left[\left(\sum_{i=0}^{\infty} \lambda^i u_i \right)^3 \right]_{\lambda=0} = 3u_3 (u_0)^2 + 6(u_0) (u_1) (u_2) + (u_1)^3 \end{aligned}$$

For $\alpha = 2, h_1(x) = x$ and $h_2(x) = 0$, we obtain the following recurrence relations,

$$\begin{aligned} u_0(x, t) &= x \\ u_{n+1}(x, t) &= -L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta [aT_x^\alpha u_{nxx}(x, t) + bu_n(x, t) + cC_n] \right] \end{aligned}$$

Consequently

$$\begin{aligned} u_0(x, t) &= x, \\ u_1(x, t) &= -L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta [au_{0xx}(x, t) + bu_0(x, t) + c(u_0)^3] \right] \\ &= -(bx + cx^3) \frac{t^\beta}{\Gamma(\beta+1)} \\ u_2(x, t) &= -L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta [au_{1xx}(x, t) + bu_1(x, t) + 3cu_1(u_0)^2] \right] \\ &= -L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta [-6acx - b(bx + cx^3) - 3c(bx + cx^3)(x)^2] \right] \\ &= [(6ac + b^2)x + 4bcx^3 + 3c^2x^5] \frac{t^{2\beta}}{\Gamma(2\beta+1)} \\ u_3(x, t) &= -L_\beta^{-1} \left[\frac{1}{s^\beta} L_\beta [au_{2xx}(x, t) + bu_2(x, t) + 3c(u_1)^2(u_0) + 3cu_2(u_0)^2] \right] \end{aligned}$$

Thus,

$$\begin{aligned} u(x, t) &= x - (bx + cx^3) \frac{t^\beta}{\Gamma(\beta+1)} \\ &+ ((6ac + b^2)x + 4bcx^3 + 3c^2x^5) \frac{t^{2\beta}}{\Gamma(2\beta+1)} + \dots \end{aligned} \quad (4.26)$$

Figure 1.4 is a graph of the solution (4.21), for $a = b = c = 1$ and α, β equal to 1.25, 1.5, 1.75, and 2.

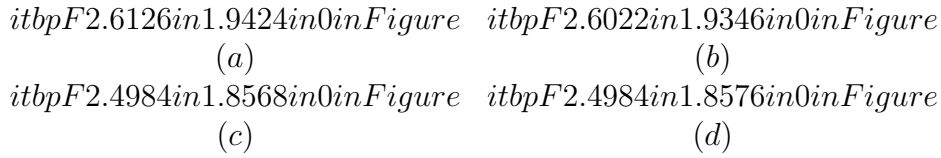


Fig1.4 Graph of Eq.(4.26) with different α, β values
 (a) $\alpha = \beta = 1.25$, (b) $\alpha = \beta = 1.5$, (c) $\alpha = \beta = 1.75$, (d) $\alpha = \beta = 2$,

5 Conclusion and Open Problem

The LDM method is applied for some nonlinear conformable fractional differential problems. It is shown that this method is a powerful device to solve not only linear problems but it is also valid for nonlinear ones.

At the end of this paper, we shall propose the following open question:

We think it is important to address a comparative study with other numerical methods for the above general conformable fractional evolution problem.

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